Reactions of Gold Atoms and Small Clusters with CO: Infrared Spectroscopic and Theoretical Characterization of Au_nCO (n = 1-5) and $Au_n(CO)_2$ (n = 1, 2) in Solid Argon

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Laser-ablated Au atoms have been co-deposited with CO molecules in solid argon to produce gold carbonyls. In addition to the previously reported Au(CO)_n (n = 1, 2) and Au₂(CO)₂ molecules, small gold cluster monocarbonyls Au_nCO (n = 2-5) are formed on sample annealing and characterized using infrared spectroscopy on the basis of the results of the isotopic substitution and CO concentration change and comparison with theoretical predictions. Of particular interest is that the mononuclear gold carbonyls, Au(CO)_n (n = 1, 2), are favored under the experimental conditions of higher CO concentration and lower laser energy, whereas the yields of the gold cluster carbonyls, Au_nCO (n = 2-5) and Au₂(CO)₂, remarkably increase with lower CO concentration and higher laser power. Density functional theory (DFT) calculations have been performed on these molecules and the corresponding small naked gold clusters. The identities of these gold carbonyls Au_nCO (n = 1-5) and Au_n(CO)₂ (n = 1, 2) are confirmed by the good agreement between the experimental and calculated vibrational frequencies, relative absorption intensities, and isotopic shifts.

I. Introduction

The interaction of CO with transition-metal and main-groupelement atoms, clusters, and surfaces plays a very important role in chemistry and chemical industry.^{1,2} Many industrial processes employ CO as a reagent and transition-metal compounds as heterogeneous catalysts and involve the intermediates of metal carbonyls. As noble-metal compounds, gold carbonyls have been the subject of a number of studies. The first carbonyl derivative of gold, Au(CO)Cl, was reported in 1925.³ Neutral gold mono- and dicarbonyls have been successfully synthesized in rare-gas matrixes by several groups.⁴ Gold carbonyl cations, $[Au(CO)_n]^+$ (n = 1, 2), have been formed in superacid solutions or in the presence of weakly coordinating counteranions⁵ and reported to be an excellent catalyst for carbonylation of olefins and alcohols.⁶ Moreover, $[Au(CO)_n]^+$ (n = 1-4) have been generated and identified by the technique of laser ablation coupled with matrix isolation infrared spectroscopy.^{4c}

Recently, it has been found that gold becomes catalytically active when deposited on select metal oxides as hemispherical ultrafine particles with diameters smaller than 5 nm.⁷ The supported Au nanoparticles exhibit remarkable catalytic activities and/or excellent selectivities in a number of reactions such as low-temperature CO oxidation and reduction of nitrogen oxides.⁷ In general, probe molecules (i.e., CO, H₂, H₂O, NO, NH₃, etc.) are used to gain insight into their chemisorption characteristics and catalytic activity. Reliable experimental data for neutral and charged clusters can be obtained, such as ionization potentials (IP), electron affinities (EA), magnetic moments, photoelectron spectra, infrared absorption spectra, polarizabilities, optical properties, and ligand adsorption capacities.8 However, it is difficult to directly determine the structure of a metal cluster due to the fact that clusters are often produced in gas-phase beams and are too small (3-50 atoms) for applying diffraction techniques.8 Fortunately, recent advances in methodology based on the technologies of pseudopotential and planewave basis sets and high-speed computers have now made it possible to obtain quantitative information on the cluster's structures as well as IP and EA, etc.^{2,8}

The technique of laser ablation coupled with matrix isolation has proven to be an efficient method to generate neutral atoms and cations of transition metals⁹ and main-group elements¹⁰ as well as electrons. In addition to a large number of mononuclear metal carbonyls, small metal cluster carbonyls, such as Fe₂- $CO_{11}^{11} Co_2 CO_{12}^{12} B_2 CO^{13}$ and $B_2 (CO)_{2}^{14} M_n CO$ (n = 2-5; M = Si, Ge),¹⁵ and Pb_nCO (n = 1-4),¹⁶ have recently been synthesized via this combined approach. Meanwhile, theoretical investigations have been carried out for Fe_nCO (n = 1-6)¹⁷ and Ni₂CO.¹⁸ In contrast with considerable studies of mononuclear gold carbonyls,³⁻⁶ however, much less is known about small gold cluster carbonyls. Here, we report the observation of the mono- and polynuclear gold carbonyls, Au_nCO (n = 1-5) and $Au_n(CO)_2$ (n = 1, 2), generated from reactions of gold atoms and small clusters with CO molecules in solid argon using infrared spectroscopy. Density functional theory (DFT) calculations have been performed to support the experimental assignments of the infrared spectra.

II. Experimental and Theoretical Methods

The experiment for laser ablation and matrix isolation infrared spectroscopy has been described previously^{16,19} and is very similar to the technique reported earlier by Andrews' group.^{9a} Briefly, the Nd:YAG laser fundamental (1064 nm, 10 Hz repetition rate with 10 ns pulse width) was focused on the rotating gold target. The laser-ablated gold atoms were co-deposited with CO in excess argon onto a CsI window cooled normally to 7 K by means of a closed-cycle helium refrigerator. Typically, 3–15 mJ/pulse laser power was used. Carbon monoxide (99.95% CO), ¹³C¹⁶O (99%, ¹⁸O < 1%), and ¹²C¹⁸O (99%) were used to prepare the CO/Ar mixtures. In general, matrix samples were deposited for 1–2 h with a typical rate of

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Figure 1. Infrared spectra in the $2160-1840 \text{ cm}^{-1}$ region from co-deposition of laser-ablated Au atoms with 0.01% CO in Ar: (a) 1 h of sample deposition at 7 K, (b) after annealing to 30 K, (c) after annealing to 34 K, (d) after 15 min of broad-band irradiation, (e) after annealing to 38 K following d, and (f) 0.01% CO + 0.002% CCl₄ after annealing to 34 K.



Figure 2. Infrared spectra in the $2160-1840 \text{ cm}^{-1}$ region for laser-ablated Au atoms co-deposited with different CO concentrations in Ar after annealing to 34 K: (a) 0.01% CO, (b) 0.08% CO, (c) 0.2% CO, and (d) 0.4% CO.

2–4 mmol/h. After sample deposition, IR spectra were recorded on a BIO-RAD FTS-6000e spectrometer at 0.5 cm⁻¹ resolution using a liquid nitrogen cooled HgCdTe (MCT) detector for the spectral range 5000–400 cm⁻¹. Samples were annealed at different temperatures and subjected to broad-band irradiation ($\lambda > 250$ nm) using a high-pressure mercury arc lamp (Ushio, 100 W).

Quantum chemical calculations were performed to predict the structures and vibrational frequencies of the observed reaction products using the GAUSSIAN 03 program.²⁰ The Becke three-parameter hybrid functional with the Lee-Yang-Parr correlation corrections (B3LYP) was used.²¹ The 6-311+G-(d) basis sets were used for C and O atoms²² and the Los Alamos ECP plus DZ for Au atoms.²³ Geometries were fully optimized and vibrational frequencies calculated with analytical second derivatives. Previous investigations have shown that the use of hybrid B3LYP along with ECP can provide reliable information for gold carbonyls, such as infrared frequencies, relative absorption intensities, and isotopic shifts.^{4c}

III. Results and Discussion

Experiments have been done with carbon monoxide concentrations ranging from 0.01% to 0.4% in argon. Typical infrared spectra for the reactions of laser-ablated gold atoms with CO molecules in excess argon in the selected regions are shown in Figures 1–5, and the absorption bands in different isotopic experiments are listed in Table 1. The stepwise annealing and photolysis behavior of the product absorptions is also shown in the figures and will be discussed below. Experiments were also done with different concentrations of CCl₄ serving as an electron scavenger in solid argon. As can be seen in Figure 1, doping with CCl₄ has no effect on these bands, suggesting that the products are neutral.

Quantum chemical calculations have been carried out for all the possible isomers of the potential product molecules. Figure



Figure 3. Infrared spectra in the $2160-1860 \text{ cm}^{-1}$ region from co-deposition of laser-ablated Au atoms with isotopic CO in Ar after annealing to 34 K: (a) 0.08% ${}^{12}\text{C}{}^{16}\text{O}$, (b) 0.04% ${}^{12}\text{C}{}^{16}\text{O} + 0.04\% {}^{13}\text{C}{}^{16}\text{O}$, and (c) 0.08% ${}^{13}\text{C}{}^{16}\text{O}$.



Figure 4. Infrared spectra in the $2160-1980 \text{ cm}^{-1}$ region for laserablated Au atoms co-deposited with isotopic CO in Ar after annealing to 34 K: (a) $0.01\% \ {}^{12}C^{16}O$, (b) $0.005\% \ {}^{12}C^{16}O + 0.005\% \ {}^{13}C^{16}O$, (c) $0.01\% \ {}^{13}C^{16}O$, (d) $0.005\% \ {}^{12}C^{16}O + 0.005\% \ {}^{12}C^{18}O$, and (e) $0.01\% \ {}^{12}C^{18}O$.

6 shows the most stable structures of the reaction products. The ground electronic states, point groups, vibrational frequencies, and intensities are listed in Table 2. Table 3 reports a comparison of the observed and calculated isotopic frequency ratios for the C-O stretching modes of the reaction products.

A. AuCO and Au(CO)₂. The sharp band at 2039.3 cm⁻¹ with a trapping site at 2037.2 cm⁻¹ (Table 1 and Figures 1 and 2) has been observed on sample deposition. The main 2039.3 cm⁻¹ band shifts to 1994.1 cm⁻¹ with ¹³C¹⁶O and to 1991.3 cm⁻¹ with ¹²C¹⁸O. The mixed ¹²C¹⁶O + ¹³C¹⁶O and ¹²C¹⁶O + ¹²C¹⁸O isotopic spectra (Figures 3 and 4) only provide the sum of pure isotopic bands, indicating that only one CO subunit is involved in this mode. The isotopic ¹²C¹⁶O/¹³C¹⁶O and ¹²C¹⁶O/¹²C¹⁸O ratios are 1.0227 and 1.0241, respectively, also indicating single CO involvement. Accordingly, the 2039.3 cm⁻¹ band is assigned to the C–O stretching mode of AuCO. The corresponding C–O stretching frequencies of AuCO in Ne, Kr, and Xe have been reported to be 2053.2, 2040.5, and 1985.5 cm⁻¹, respectively.^{4a,c}

We also note that the 1916.3, 1925.4, 1928.2, and 1934.1 cm^{-1} bands increase together on sample annealing. Interestingly,



Figure 5. Infrared spectra in the 1940–1800 cm⁻¹ region for laserablated Au atoms co-deposited with isotopic CO in Ar after annealing to 34 K: (a) $0.01\% \ ^{12}C^{16}O$, (b) $0.005\% \ ^{12}C^{16}O + 0.005\% \ ^{13}C^{16}O$, (c) $0.01\% \ ^{13}C^{16}O$, (d) $0.005\% \ ^{12}C^{16}O + 0.005\% \ ^{12}C^{18}O$, and (e) $0.01\% \ ^{12}C^{18}O$.

the yields of these bands relative to the 2039.3 cm⁻¹ band for AuCO sharply increase with higher CO concentration (i.e., 0.4%) and lower laser power (5 mJ/pulse) (Figure 2). As can be seen in Figure 3, four sets of triplet bands have been observed at 1916.3/1890.5/1871.7, 1925.4/1898.1/1879.4, 1928.2/1901.9/ 1883.3, and 1934.1/1907.7/1889.2 cm⁻¹ in the mixed ¹²C¹⁶O + ¹³C¹⁶O isotopic spectra (Table 1). Similar isotopic spectra in the ¹²C¹⁶O + ¹²C¹⁸O experiments have also been obtained. These bands are assigned to the antisymmetric C–O stretching modes of the Au(CO)₂ molecule in different matrix sites, in agreement with previous studies.^{4a,c} Weak bands have been observed for the symmetric CO stretching modes for Au(¹²CO)(¹³CO) and Au(C¹⁶O)(C¹⁸O) (Table 1).

Density functional theory (DFT) calculations lend strong support for the assignments. AuCO is predicted to have a bent geometry (Figure 6) with a ²A' ground state (Table 2), which lies 4.51 kcal/mol lower than the linear structure. The CO binding energy in AuCO is 8.19 kcal/mol, which is in excellent agreement with the previous theoretical prediction (8.2 kcal/mol).^{4c} The calculated C–O stretching frequency of the AuCO species is 2065.9 cm⁻¹ (Table 2), which should be multiplied by 0.987 to fit the observed frequency. The calculated ¹²C¹⁶O/

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Figure 6. Optimized structures (bond lengths in Angstroms, bond angles in degrees) of the most stable isomers for the reaction products calculated at the B3LYP/6-311+G(d)-LANL2DZ level.

TABLE 1: Infrared Absorptions (cm $^{-1}$) Observed from the Reactions of Laser-Ablated Gold Atoms with CO in Excess Argon at 7 K

¹² C ¹⁶ O	$^{13}C^{16}O$	$^{12}C^{18}O$	${}^{12}C^{16}O + {}^{13}C^{16}O$	${}^{12}C^{16}O + {}^{12}C^{18}O$	<i>R</i> (12/13)	<i>R</i> (16/18)	assignment
2131.9	2083.0	2084.1	2131.9, 2083.1	2131.9, 2084.2	1.0235	1.0229	Au ₂ CO
2115.9	2067.2	2068.8	2115.9, 2067.2	2115.9, 2068.8	1.0236	1.0228	Au ₄ CO
2115.0	2066.4	2068.1	2115.0, 2066.4	2115.0, 2068.1	1.0235	1.0227	Au ₄ CO site
2111.0	2064.5	2064.1	2111.8, 2073.0, 2064.5, 2123.4 (sym)		1.0225	1.0227	Au ₂ (CO) ₂ site
2110.2	2062.8	2061.2	2110.2, 2069.5, 2062.8, 2120.0 (sym)	2110.2, 2068.6, 2061.2, 2120.2 (sym)	1.0228	1.0237	$Au_2(CO)_2$
2075.4	2028.6	2027.9	2075.4, 2028.6	2075.4, 2028.0	1.0231	1.0234	Au ₃ CO
2039.3	1994.1	1991.3	2039.5, 1994.3	2039.4, 1991.2	1.0227	1.0241	AuCO
2037.2	1992.2	1989.1	2037.3, 1992.2	2037.3, 1989.1	1.0226	1.0241	AuCO site
1934.1	1889.1	1892.6	1934.1, 1907.7, 1889.2, 2056.0 (sym)	1934.1, 1910.2, 1892.6, 2071.9 (sym)	1.0238	1.0219	$Au(CO)_2$ site
1928.2	1883.3	1886.9	1928.2, 1901.9, 1883.3, 2051.8 (sym)	1928.2, 1904.5, 1886.9, 2062.7 (sym)	1.0238	1.0219	$Au(CO)_2$ site
1925.4	1879.4	1884.4	1925.4, 1898.1, 1879.4, 2049.6 (sym)	1925.4, 1900.9, 1884.4	1.0238	1.0218	$Au(CO)_2$ site
1916.3	1871.7	1875.5	1916.3, 1890.5, 1871.7, 2047.0 (sym)	1916.3, 1893.1, 1875.5, 2059.4 (sym)	1.0238	1.0218	$Au(CO)_2$
1854.8	1814.2	1811.8	1854.8, 1814.2	1854.8, 1811.8	1.0224	1.0237	Au ₅ CO site
1852.9	1812.1	1809.6	1852.9, 1812.0	1852.9, 1809.6	1.0225	1.0239	Au ₅ CO

TABLE 2: Ground Electronic States, Point Groups, Vibrational Frequencies (cm $^{-1}$), and Intensities (km/mol) of the Reaction Products Calculated at the B3LYP/6-311+G(d)-LANL2DZ Level

species	electronic state	point group	frequency (intensity, mode)
AuCO	² A'	Cs	2065.9 (832, A'), 329.5 (5, A'), 210.7 (15, A')
$Au(CO)_2$	${}^{2}A_{1}$	$C_{2\mathrm{v}}$	2103.4 (4, A ₁), 1980.5 (3957, B ₂), 389.1 (58, B ₂), 377.0 (0.1, A ₁), 296.5 (79, B ₂), 257.6 (3, A ₁),
			$257.6 (0, A_2), 132.7 (4, B_1), 53.0 (0.4, A_1)$
$Au_2(CO)_2$	$^{1}\Sigma_{g}^{+}$	$D_{\infty h}$	2186.7 (0, $\sigma_{\rm g}$), 2172.0 (1118, $\sigma_{\rm u}$), 319.7 (0 × 2, π), 305.7 (0, $\sigma_{\rm g}$), 260.6 (16 × 2, π), 258.9 (0, $\sigma_{\rm u}$),
			$162.7 (0, \sigma_{g}), 26.0 (0 \times 2, \pi), 21.1 (0 \times 2, \pi)$
Au ₂ CO	${}^{1}A'$	C_s	2192.8 (588, Å'), 368.4 (0.1, A'), 335.9 (9.0, A'), 282.5 (3, A''), 167.4 (4, A'), 47.5 (0, A')
Au ₃ CO	${}^{2}\mathbf{B}_{2}$	C_{2v}	2169.2 (1226, A ₁), 391.4 (6, A ₁), 342.4 (5, B ₁), 328.6 (1, B ₂), 147.6 (3, A ₁), 91.0 (0.3, A ₁),
			81.5 (2, B ₂), 52.7 (0, B ₁), 30.6(0.2, B ₂)
Au ₄ CO	${}^{1}A'$	Cs	2198.0 (819, A'), 380.6 (7, A'), 340.4 (1, A"), 335.8 (1, A'), 178.1 (5, A'), 139.7 (2, A'), 90.8 (2, A'),
			75.8 (0.3, A'), 51.2 (0, A''), 36.0 (0, A'), 33.9 (0.1, A''), 19.6 (0.4, A')
Au ₅ CO	${}^{2}\mathbf{B}_{2}$	C_{2v}	1923.7 (663, A ₁), 376.1 (1, B ₁), 369.5 (21, B ₂), 296.6 (1, A ₁), 188.9 (4, B ₂), 157.0 (5, B ₂), 131.2 (1, A ₁),
			95.8 (0, B ₂), 93.8(0, A ₁), 79.2 (0.3, A ₁), 50.8 (0, A ₁), 49.1 (0.2, B ₂), 46.3 (0.3, B ₁), 33.6 (0.2, B ₁),
			$30.1(0, A_2)$

¹³C¹⁶O and ¹²C¹⁶O/¹²C¹⁸O isotopic frequency ratios of 1.0225 and 1.0252 (Table 3) are consistent with the experimental values, 1.0227 and 1.0241, respectively. Similarly, the good agreement between the experimental and calculated vibrational frequencies, relative absorption intensities, and isotopic shifts confirms the identification of the Au(CO)₂ molecule. Briefly, Au(CO)₂ is predicted to have a ²A₁ ground state with $C_{2\nu}$ symmetry and \angle CAuC = 176.9° (Table 2 and Figure 6), consistent with the previous report.^{4c}

B. $Au_2(CO)_2$. The absorption at 2110.2 with a trapping site at 2111.0 cm⁻¹ appears on sample deposition, observably increases on annealing, decreases on broad-band irradiation, and recovers on further annealing, as shown in Figure 1. The 2110.2

TABLE 3: Comparison of the Observed and CalculatedIsotopic Frequency Ratios of the Reaction Products

		$^{12}C^{16}O$	$^{12}C^{16}O/^{13}C^{16}O$		$^{12}\mathrm{C}^{16}\mathrm{O}/^{12}\mathrm{C}^{18}\mathrm{O}$	
molecule	mode	obsd	calcd	obsd	calcd	
AuCO	C–O str.	1.0227	1.0225	1.0241	1.0252	
$Au(CO)_2$	C-O asym-str.	1.0238	1.0235	1.0218	1.0238	
$Au_2(CO)_2$	C-O asym-str.	1.0230	1.0230	1.0238	1.0244	
Au ₂ CO	C–O str.	1.0235	1.0235	1.0229	1.0237	
Au ₃ CO	C-O str.	1.0231	1.0236	1.0234	1.0235	
Au ₄ CO	C-O str.	1.0236	1.0236	1.0228	1.0236	
Au ₅ CO	C–O str.	1.0225	1.0228	1.0239	1.0247	

 cm^{-1} band shifts to 2064.5 cm^{-1} with ${}^{13}C^{16}O$ and to 2061.2 cm^{-1} with ${}^{12}C^{18}O$ (Table 1, Figures 3 and 4). In the mixed

 ${}^{12}C^{16}O + {}^{13}C^{16}O$ experiment, a triplet at 2110.2, 2069.5, and 2062.8 cm⁻¹ together with an associated band at 2120.0 cm⁻¹ is observed. A similar isotopic splitting feature is obtained in the mixed ${}^{12}C^{16}O + {}^{12}C^{18}O$ isotopic spectra, indicating that two CO subunits are involved. As can be seen in Figure 2, the 2110.2 cm⁻¹ band sharply increases with lower CO concentration (0.01%) and higher laser power (12 mJ/pulse) (Figure 2a), whereas the reverse is true for the relative yields of $Au(CO)_2$ bands. In the 0.01% CO and 12 mJ/pulse experiment the intensity ratio of the 2110.2, 2039.3, and 1916.4 cm^{-1} bands is 0.080:0.022:0.031, while this ratio changes to 0.019:0.052:0.550 in the 0.08% CO and 5 mJ/pulse experiment and to 0.009:0.139: 0.280 in the 0.2% CO and 5 mJ/pulse experiment, suggesting that the molecule with the 2110.2 cm^{-1} band involves more Au atoms than AuCO and Au(CO)₂. This band is assigned to the antisymmetric C-O stretching mode of Au₂(CO)₂. The bands observed at 2120.0 cm⁻¹ in the mixed ${}^{12}C^{16}O + {}^{13}C^{16}O$ experiment and 2120.2 cm⁻¹ in the ${}^{12}C^{16}O + {}^{12}C^{18}O$ experiment are due to the symmetric C–O stretching modes of $Au_2(^{12}CO)$ - (^{13}CO) and Au₂(C¹⁶O)(C¹⁸O), respectively. The antisymmetric C-O stretching vibrations at 2110.2 and 2111.0 cm⁻¹ are compatible with the previous neon matrix studies with 6.4 and 7.0 cm⁻¹ blue shifts.^{4c}

The assignment is strongly supported by the present DFT calculations, which predict this Au₂(CO)₂ molecule to have a linear geometry with a ${}^{1}\Sigma_{g}^{+}$ ground electronic state (Table 2 and Figure 6), in accord with the previous report,^{4c} whereas nonlinear structures are observed for AuCO,^{4c} Au(CO)₂,^{4c} and Au₂CO (Figure 6). The antisymmetric C–O stretching frequency is calculated to be 2172.0 cm⁻¹, which requires a 0.972 scale factor. The calculated ${}^{12}C^{16}O/{}^{13}C^{16}O$ and ${}^{12}C^{16}O/{}^{12}C^{18}O$ isotopic frequency ratios of 1.0230 and 1.0244 (Table 3) are again in excellent agreement with the experimental values, 1.0230 and 1.0238, respectively.

C. Au_nCO (n = 2-5). The present infrared spectra provide evidence for the formation of small gold cluster carbonyls in the excess argon matrixes. For instance, weak bands at 2131.9, 2115.9, 2075.4, and 1852.9 cm⁻¹ appear on sample deposition, which observably increase on sample annealing. The experimental conditions of lower CO concentration and higher laser power favors the formation of species with these bands (Figure 2). These bands have some trapping site absorptions (Table 1 and Figure 1), and we will focus on the main bands.

The 2131.9, 2115.9, 2075.4, and 1852.9 cm⁻¹ absorptions shift to 2083.0, 2067.2, 2028.6, and 1812.1 cm⁻¹ with ${}^{13}C^{16}O$ and to 2084.1, 2068.8, 2027.9, and 1809.6 cm^{-1} with ${}^{12}C^{18}O$, respectively. In the mixed ${}^{12}C^{16}O + {}^{13}C^{16}O$ and ${}^{12}C^{16}O +$ ¹²C¹⁸O experiments, only pure isotopic counterparts are observed. The isotopic ratios (12C16O/13C16O:1.0235, 1.0236, 1.0231, and 1.0225; ¹²C¹⁶O/¹²C¹⁸O: 1.0229, 1.0228, 1.0234, and 1.0239, respectively) and the mixed isotopic characteristic (Figures 4 and 5) indicate that only one CO subunit is involved in each mode. On the basis of their different isotopic frequency shifts and different annealing and photolysis behavior, each of these bands is attributed to one of the gold cluster carbonyls, Au_nCO with $n \ge 2$. The 2131.9, 2115.9, and 2075.4 cm⁻¹ absorptions lie in the region expected for terminal-bonded C-O stretching vibrations, whereas the 1852.9 cm^{-1} band is due to a bridge-bonded C-O stretching vibration. The series experiments with different CO concentrations and laser energies show that the appearance of the 2131.9 cm⁻¹ band is prior to the other three bands, which is assigned to the C-O stretching mode of Au₂CO. It is expected that a high Au/CO ratio favors the formation of higher clusters during annealing, as discussed in

detail later (vide infra). The observed annealing behavior and the Au/CO ratio dependence, along with the DFT predictions, indicate that the bands at 2075.4, 2115.9, and 1852.9 cm⁻¹ can be assigned to the C–O stretching vibrations of Au₃CO, Au₄-CO, and Au₅CO, respectively.

It can be found from Tables 1-3 that the calculated frequencies (only 0.03%, 0.05%, 0.04%, and 0.04% higher, respectively) as well as the isotopic ratios are in excellent agreement with the experimental values. As shown in Figure 6, among the most stable gold cluster monocarbonyls, a terminal CO is found in Au₂CO, Au₃CO, or Au₄CO whereas a bridging CO in Au₅CO.

The present optimization results show that the naked Au₂ cluster has a ${}^{1}\Sigma_{g}{}^{+}$ ground state with the Au–Au bond of 2.574 Å, in accord with the experimental values (~2.47 Å)²⁴ and the previous theoretical calculations (~2.60 Å).²⁵ The Au–Au stretching of the Au₂ cluster is IR inactive, whereas the vibrational frequency of Au₂ was observed at 191 cm⁻¹ in a study of emission and laser excitation spectra.²⁴ The most stable structure of the Au₂CO molecule is predicted to be asymmetric with a ¹A' ground state (Table 2 and Figure 6), which lies 27.77 kcal/mol lower than the bridge-bonded Au₂CO structure.

The naked Au₃ cluster has a ²B₂ ground state with $C_{2\nu}$ symmetry, exhibiting an open obtuse angle structure with an apex angle of 64.8°, which is in accordance with the previous studies.^{25a,b} Our DFT calculations predict that Au₃CO has a ²B₂ ground state with an Au–C bond of 1.949 Å. In Au₃CO, the apex angle (62.2°) is 2.6° smaller than that of Au₃ at the same DFT level. The bridge-bonded Au₃CO structure lies 23.31 kcal/ mol higher than the most stable terminal-bonded one.

For Au₄, a planar rhombus arrangement with D_{2h} symmetry in the ${}^{1}A_{g}$ ground state is the most stable structure, in agreement with the previous studies.^{25a,b} The most stable structure of Au₄-CO is predicted to have a planar geometry with C_{s} symmetry, as shown in Figure 6. The CO molecule is terminal bonded to one of the apex Au atoms of the triangle Au₃ ring. The \angle CAuAu is calculated to be 162.4°. In the less stable Au₄CO, CO is terminal bonded to one of the apex Au atoms of a near rhombus Au₄. It is found that the Au₄CO with a near pyramid structure is 157.10 kcal/mol higher than the most stable one.

The naked Au₅ cluster is predicted to have a ${}^{2}A_{1}$ ground state with $C_{2\nu}$ symmetry, in line with previous theoretical calculations.^{25a,b} The most stable Au₅CO structure has a ${}^{2}B_{2}$ ground state with $C_{2\nu}$ symmetry, and CO is bridge bonded to two Au atoms of Au₅ (Figure 6). In contrast, Si₅CO was predicted to have a nonplanar structure with $C_{2\nu}$ symmetry in which CO bridges two of the equatorial silicon atoms of Si₅ with a trigonal bipyramid structure.^{15a}

In addition, the calculated net charges of CO (not listed here) show that the partial electron is transferred from Au_n (n = 1-5) to CO, where Au_n behaves as a donor while CO behaves as an acceptor in bonding between Au_n and CO. Then, the C–O bond gains some activation as expected.

D. Reaction Mechanisms. With low CO concentrations and high laser energies the laser-ablated gold atoms react with CO molecules in the excess argon matrixes to produce the small gold cluster carbonyls, Au_nCO (n = 2-5) and Au₂(CO)₂, in addition to Au(CO)_n (n = 1, 2) (Figure 1).

Under the present experimental conditions, gold atoms are the predominant species produced by laser ablation of the gold target. Only very small absorptions of the primary Au_nCO (n = 2-5) species are observed upon deposition, and these species appear and increase significantly upon sample annealing to high temperatures (30–34 K) in experiments with low CO concen-

TABLE 4: Energetics for Possible Reactions of Gold Atoms and Small Clusters with CO Calculated at the B3LYP/ 6-311+G(d)-LANL2DZ Level

no.	reaction	reaction energy ^a (kcal/mol)
1	Au $({}^{2}S_{1/2})$ + Au $({}^{2}S_{1/2}) \rightarrow$ Au ₂ $({}^{1}\Sigma_{\sigma}^{+})$	-43.20
2	$\operatorname{Au}_2({}^{1}\Sigma_{\sigma}^{+}) + \operatorname{Au}({}^{2}S_{1/2}) \rightarrow \operatorname{Au}_3({}^{2}B_2)$	-20.27
3	$\operatorname{Au}({}^{2}\operatorname{S}_{1/2}) + \operatorname{Au}_{3}({}^{2}\operatorname{B}_{2}) \rightarrow \operatorname{Au}_{4}({}^{1}\operatorname{A}_{2})$	-44.04
4	$\operatorname{Au}_2({}^{1}\Sigma_g^{+}) + \operatorname{Au}_2({}^{1}\Sigma_g^{+}) \rightarrow \operatorname{Au}_4({}^{1}A_g)$	-21.11
5	Au $(^{2}S_{1/2})$ + Au ₄ $(^{1}A_{g}) \rightarrow$ Au ₅ $(^{2}A_{1})$	-40.22
6	$\operatorname{Au}_2({}^{1}\Sigma_g^{+}) + \operatorname{Au}_3({}^{2}B_2) \rightarrow \operatorname{Au}_5({}^{2}A_1)$	-41.06
7	Au $({}^{2}S_{1/2})$ + CO $({}^{1}\Sigma^{+}) \rightarrow$ AuCO $({}^{2}A')$	-8.19
8	AuCO $(^{2}A')$ + CO $(^{1}\Sigma^{+}) \rightarrow$ Au(CO) ₂ $(^{2}A_{1})$	-13.04
9	$\operatorname{Au}_2({}^{1}\Sigma_g^{+}) + \operatorname{CO}({}^{1}\Sigma^{+}) \rightarrow \operatorname{Au}_2\operatorname{CO}({}^{1}A')$	-25.56
10	AuCO $({}^{2}A')$ + Au $({}^{2}S_{1/2}) \rightarrow$ Au ₂ CO $({}^{1}A')$	-60.56
11	AuCO $(^{2}A')$ + AuCO $(^{2}A')$ \rightarrow Au ₂ (CO) ₂ $(^{1}\Sigma_{g}^{+})$	-65.09
12	$Au(CO)_2(^2A_1) + Au(^2S_{1/2}) \rightarrow Au_2(CO)_2(^1\Sigma_g^{+})$	-60.25
13	$\operatorname{Au_2CO}({}^{1}\mathrm{A'}) + \operatorname{CO}({}^{1}\Sigma^+) \rightarrow \operatorname{Au_2(CO)_2}({}^{1}\Sigma_{g}^+)$	-12.72
14	$\operatorname{Au}_{3}(^{2}\mathrm{B}_{2}) + \operatorname{CO}(^{1}\Sigma^{+}) \rightarrow \operatorname{Au}_{3}\operatorname{CO}(^{2}\mathrm{B}_{2})$	-26.27
15	AuCO $(^{2}A')$ + Au ₂ $(^{1}\Sigma_{g}^{+}) \rightarrow$ Au ₃ CO $(^{2}B_{2})$	-38.35
16	$Au_2CO(^1A') + Au(^2S_{1/2}) \rightarrow Au_3CO(^2B_2)$	-20.99
17	$\operatorname{Au}_4({}^{1}A_g) + \operatorname{CO}({}^{1}\Sigma^+) \rightarrow \operatorname{Au}_4\operatorname{CO}({}^{1}A')$	-29.40
18	AuCO $({}^{2}A')$ + Au ₃ $({}^{2}B_{2}) \rightarrow$ Au ₄ CO $({}^{1}A')$	-65.25
19	$Au_2CO({}^{1}A') + Au_2({}^{1}\Sigma_g^+) \rightarrow Au_4CO({}^{1}A')$	-24.96
20	$Au_3CO(^2B_2) + Au(^2S_{1/2}) \rightarrow Au_4CO(^1A')$	-47.17
21	$\operatorname{Au}_{5}(^{2}A_{1}) + \operatorname{CO}(^{1}\Sigma^{+}) \rightarrow \operatorname{Au}_{5}\operatorname{CO}(^{2}B_{2})$	-13.65
22	AuCO $(^{2}A')$ + Au ₄ $(^{1}A_{g}) \rightarrow$ Au ₅ CO $(^{2}B_{2})$	-45.67
23	$Au_2CO(^1A') + Au_3(^2B_2) \rightarrow Au_5CO(^2B_2)$	-29.15
24	$\operatorname{Au_3CO}(^2B_2) + \operatorname{Au_2}(^1\Sigma_g^+) \rightarrow \operatorname{Au_5CO}(^2B_2)$	-28.43
25	$\operatorname{Au_4CO}(^1A') + \operatorname{Au}(^2S_{1/2}) \rightarrow \operatorname{Au_5CO}(^2B_2)$	-24.46

^{*a*} A negative value of energy denotes that the reaction is exothermic.

trations (see Figure 1), which indicates that these small gold clusters are mainly formed in solid argon upon annealing but not during the laser ablation process. This means that the higher laser power leads to neither generating more clusters in the gas phase nor significant annealing of the matrix during deposition and so increasing of the yield of Au_n species. It seems that a higher laser power accounts for the generation of a higher concentration Au in the matrix, corresponding to a higher Au/CO ratio.

The energetic analysis for possible reactions of gold atoms and small clusters with CO molecules has been performed at the B3LYP/6-311+G(d)-LANL2DZ level. As can be seen in Table 4, all the possible reactions of gold atoms and small clusters with CO are exothermic (-8.19 to -65.25 kcal/mol). This implies that the diffusion during annealing may make a species, i.e., a gold atom or cluster, a CO molecule, or a gold carbonyl intermediate, in the matrix readily react with its nearest neighbor, and consequently, the final product distribution will mainly depend on the Au/CO ratio, in agreement with our observations that low CO concentration and high laser power favor high nuclearity of cluster molecules (Figure 2). In contrast, at high CO concentration and low laser power, corresponding to a low Au/CO ratio, mononuclear gold polycarbonyl species are the predominant products.

IV. Conclusions

Reactions of laser-ablated Au atoms with CO molecules in solid argon have been studied using matrix-isolation infrared spectroscopy. In addition to the previously reported Au(CO)_n (n = 1, 2) and Au₂(CO)₂ molecules, small gold cluster monocarbonyls Au_nCO (n = 2-5) are observed in the present infrared spectra. On the basis of the isotopic substitution, change of CO concentration, and comparison with density functional calculations, the absorptions at 2131.9, 2075.4, 2115.9, and 1852.9 cm⁻¹ are assigned to Au₂CO, Au₃CO, Au₄CO, and Au₅-

CO, respectively. Among the most stable gold cluster monocarbonyls, a terminal CO is found in Au₂CO, Au₃CO, or Au₄CO whereas a bridging CO in Au₅CO. The experimental results of Au_nCO (n = 1-5) and Au_n(CO)₂ (n = 1, 2) are in good agreement with the prediction of density functional theory calculations.

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